

Perturbation of the loop measure

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Abstract

The loop measure is associated with a Markov generator. We compute the variation of the loop measure induced by an infinitesimal variation of the generator affecting the killing rates or the jumping rates.

1 Introduction

Professor Fukushima's contribution to probabilistic potential theory was seminal. His book on Dirichlet spaces [7] was the first complete exposition in which the functional analytic, potential theoretic, and probabilistic aspects of the theory were considered jointly, as different aspects of the same mathematical object. In particular, transformations of the energy form, such as restriction to an open set, trace on a closed set, change of the equilibrium (or killing) measure, change of reference measure, etc... have probabilistic counterparts. In the second chapter, the possibility of superposing different Dirichlet forms was recognized. The present paper follows this line of research. But we will focus on Markov loops and bridges rather than Markov paths. We will present them in the next section. Let us stress however that the existence of loop or bridge measures seems to require more than the assumption of a Markov process defined up to polar sets, which is the basic assumption for a Markov process associated with a Dirichlet form. The existence of a Green function seems to be required. Our purpose is to compute the variation of the loop measure induced by an infinitesimal variation of the generator. This variation may in particular affect the killing rates or the jumping rates. In the case of symmetric continuous time Markov chains, the question has been addressed in chapter 6 of [9]. The results are formally close to formulas used in conformal field theory for operator insertions (see for example [8]). We try here to extend them to a more general situation, to show in particular that the bridge measures can be derived from the loop measure.

2 Background on loop measures

We begin by introducing loop measures for Borel right processes (such as Feller processes) on a rather general state space S , which we assume to be locally compact with a countable base. Let $X = (\Omega, \mathcal{F}_t, X_t, \theta_t, P^x)$ be a transient Borel right process [12] with cadlag paths (such as a standard Markov process [1]) with state space S and jointly measurable transition densities $p_t(x, y)$ with respect to some σ -finite measure m on S . We assume that the potential densities

$$u(x, y) = \int_0^\infty p_t(x, y) dt$$

are finite off the diagonal, but allow them to be infinite on the diagonal. We do not require that $p_t(x, y)$ is symmetric.

We assume furthermore that $0 < p_t(x, y) < \infty$ for all $0 < t < \infty$ and $x, y \in S$, and that there exists another Borel right process \hat{X} in duality with X (Cf [1]), relative to the measure m , so that its transition probabilities $\hat{P}_t(x, dy) = p_t(y, x) m(dy)$. These conditions allow us to use material on bridge measures in [6]. In particular, for all $0 < t < \infty$ and $x, y \in S$, there exists a finite measure $Q_t^{x,y}$ on \mathcal{F}_{t-} , of total mass $p_t(x, y)$, such that

$$Q_t^{x,y} (1_{\{\zeta > s\}} F_s) = P^x (F_s p_{t-s}(X_s, y)), \quad (2.1)$$

for all $F_s \in \mathcal{F}_s$ with $s < t$. (We use the letter Q for measures which are not necessarily of mass 1, and reserve the letter P for probability measures.) $Q_t^{x,y}$ should be thought of as a measure for paths which begin at x and end at y at time t . When normalized, this gives the bridge measure $P_t^{x,y}$ of [6].

We use the canonical representation of X in which Ω is the set of cadlag paths ω in $S_\Delta = S \cup \Delta$ with $\Delta \notin S$, and is such that $\omega(t) = \Delta$ for all $t \geq \zeta = \inf\{t > 0 \mid \omega(t) = \Delta\}$. Set $X_t(\omega) = \omega(t)$. We define a σ -finite measure μ on (Ω, \mathcal{F}) by

$$\int F d\mu = \int_0^\infty \frac{1}{t} \int Q_t^{x,x} (F \circ k_t) dm(x) dt \quad (2.2)$$

for all \mathcal{F} measurable functions F on Ω . Here k_t is the killing operator defined by $k_t \omega(s) = \omega(s)$ if $s < t$ and $k_t \omega(s) = \Delta$ if $s \geq t$, so that $k_t^{-1} \mathcal{F} \subset \mathcal{F}_{t-}$. As usual, if F is a function, we often write $\mu(F)$ for $\int F d\mu$. μ is σ -finite as any set of loops in a compact set with lifetime bounded away from zero and infinity has finite measure.

We call μ the loop measure of X because, when X has continuous paths, μ is concentrated on the set of continuous loops with a distinguished starting

point (since $Q_t^{x,x}$ is carried by loops starting at x). Moreover, it is shift invariant. More precisely let ρ_u denote the loop rotation defined by

$$\rho_u \omega(s) = \begin{cases} \omega(s + u \bmod \zeta(\omega)), & \text{if } 0 \leq s < \zeta(\omega) \\ \Delta, & \text{otherwise.} \end{cases}$$

Here, for two positive numbers a, b we define $a \bmod b = a - mb$ as the unique positive integer m such that $0 \leq a - mb < b$. μ is invariant under ρ_u , for any u . We let \mathcal{F}_ρ denote the σ -algebra of \mathcal{F} measurable functions F on Ω which are invariant under ρ , that is $F \circ \rho_u = F$ for all $u \geq 0$. Loop functionals of interest are mostly \mathcal{F}_ρ -measurable. Recall that Poisson processes of intensity μ appear naturally as produced by loop erasure in the construction of random spanning trees through the Wilson algorithm (see chapter 8 in [9]). Although the definition of μ in (2.2), especially the $\frac{1}{t}$, may look forbidding, μ often has a nice form when applied to specific functions in \mathcal{F}_ρ . A particular function in \mathcal{F}_ρ is given by

$$\phi(f) = \int_0^\infty f(X_t) dt, \quad (2.3)$$

where f is any measurable function on S . If f_j , $j = 1, \dots, k \geq 2$ are non-negative functions on S , then by [11, Lemma 2.1]

$$\begin{aligned} & \mu \left(\prod_{j=1}^k \phi(f_j) \right) \\ &= \sum_{\pi \in \mathcal{P}_k^\odot} \int u(y_{\pi(1)}, y_{\pi(2)}) \cdots u(y_{\pi(k-1)}, y_{\pi(k)}) u(y_{\pi(k)}, y_{\pi(1)}) \prod_{j=1}^k f_j(y_j) dy_j \end{aligned} \quad (2.4)$$

where \mathcal{P}_k^\odot denotes the set of permutations of $[1, k]$ on the circle. (For example, $(1, 3, 2)$, $(2, 1, 3)$ and $(3, 2, 1)$ are considered to be one permutation $\pi \in \mathcal{P}_3^\odot$.) We note however that in general when u is infinite on the diagonal

$$\mu(\phi(f_j)) = \infty.$$

For $k \geq 2$, the integral (2.4) can be finite if the f_i satisfy certain integrability conditions: see the beginning of section 3. Consider more generally the multiple integral

$$\sum_{\pi \in \mathcal{T}_k^\odot} \int_{0 \leq r_1 \leq \dots \leq r_k \leq t} f_{\pi(1)}(X_{r_1}) \cdots f_{\pi(k)}(X_{r_k}). \quad (2.5)$$

where \mathcal{T}_k^\odot denotes the set of translations π of $[1, k]$ which are cyclic mod k , that is, for some i , $\pi(j) = j + i \pmod k$, for all $j = 1, \dots, k$.

Finite sums of multiple integrals such as these form an algebra (see exercise 11, p.21 in [9]) which generates \mathcal{F}_ρ , [3].

Finally, let $a(x)$ be a bounded, strictly positive function on S . Define the time changed process $Y_t = X_{\tau(t)}$ where $\tau(t) = \int_0^t a(Y_s) ds$ is the inverse of the CAF $A_t = \int_0^t 1/a(X_s) ds$. It satisfies the duality assumption relative to the measure $a \cdot m$. It then follows as in [5, section 7.3] that if $u_X(x, y), u_Y(x, y)$ denote the potential densities of X, Y respectively with respect to $m, a \cdot m$ respectively, then

$$u_Y(x, y) = u_X(x, y)/a(y). \quad (2.6)$$

It follows that μ_Y is the image of μ_X by the time change.

3 Multiple CAF's and perturbation of loop measures

Let $\mathcal{M}(S)$ be the set of finite signed Radon measures on $\mathcal{B}(S)$. We say that a norm $\|\cdot\|$ on $\mathcal{M}(S)$ is a **proper norm** with respect to a kernel u if for all $n \geq 2$ and ν_1, \dots, ν_n in $\mathcal{M}(S)$

$$\left| \int \prod_{j=1}^n u(x_j, x_{j+1}) \prod_{i=1}^n d\nu_j(x_i) \right| \leq C^n \prod_{j=1}^n \|\nu_j\|, \quad (3.1)$$

(with $x_{n+1} = x_1$) for some universal constant $C < \infty$. In Section 6 of [10] we present several examples of proper norms which depend upon various hypotheses about the kernel u .

In particular, the following norm is related to the square root of the generator of X , which defines the Dirichlet space in the m -symmetric case:

$$\|\nu\|_w := \left(\iint \left(\int w(x, y) w(y, z) d\nu(y) \right)^2 dm(x) dm(z) \right)^{1/2}, \quad (3.2)$$

where

$$w(x, y) = \int_0^\infty \frac{p_s(x, y)}{\sqrt{\pi s}} ds. \quad (3.3)$$

To see that $\|\nu\|_w$ is a proper norm we first note that

$$u(x, z) = \int w(x, y) w(y, z) dm(y) \quad (3.4)$$

(It is interesting to note that w is the potential density of the process X_{T_t} where T_t is the stable subordinator of index $1/2$. In operator notation (3.4) says that $W^2 = U$ where W and U are operators with kernels w and u respectively.)

Using (3.4)

$$\begin{aligned} \prod_{j=1}^n u(z_j, z_{j+1}) &= \prod_{j=1}^n \int w(z_j, \lambda_j) w(\lambda_j, z_{j+1}) dm(\lambda_j) \\ &= \prod_{j=1}^n \int w(z_j, \lambda_j) w(\lambda_{j-1}, z_j) dm(\lambda_j) \end{aligned} \quad (3.5)$$

in which $z_{n+1} = z_1$ and $\lambda_0 = \lambda_n$. It follows from this that

$$\begin{aligned} &\left| \int \prod_{j=1}^n u(z_j, z_{j+1}) \prod_{j=1}^n d\nu_j(z_j) \right| \\ &= \left| \int \prod_{j=1}^n \left(\int w(z_j, \lambda_j) w(\lambda_{j-1}, z_j) d\nu_j(z_j) \right) \prod_{j=1}^n dm(\lambda_j) \right| \\ &\leq \prod_{j=1}^n \left(\iint \left(\int w(z_j, s) w(t, z_j) d\nu_j(z_j) \right)^2 dm(s) dm(t) \right)^{1/2} \leq \prod_{j=1}^n \|\nu_j\|_w, \end{aligned} \quad (3.6)$$

where, for the first inequality, we use repeatedly the Cauchy-Schwarz inequality.

Lastly, set

$$M_\nu(x, z) = \int w(x, y) w(y, z) d\nu(y). \quad (3.7)$$

Since $\|\nu\|_w$ is the L^2 norm of M_ν , and $M_{\nu+\nu'} = M_\nu + M_{\nu'}$, we see that $\|\nu\|_w$ is a norm. (This can also be viewed as the Hilbert-Schmidt norm of the operator defined by the kernel M_ν .) \square

We denote by \mathcal{R}^+ the set of positive bounded Revuz measures ν on S that are associated with X . This is explained in detail in Section 2.1 of [10]. We use L_t^ν to denote the CAF with Revuz measure ν .

Let $\|\cdot\|$ be a proper norm on $\mathcal{M}(S)$ with respect to the kernel u . Set

$$\mathcal{M}_{\|\cdot\|}^+ = \{\text{positive } \nu \in \mathcal{M}(S) \mid \|\nu\| < \infty\}, \quad (3.8)$$

and

$$\mathcal{R}_{\|\cdot\|}^+ = \mathcal{R}^+ \cap \mathcal{M}_{\|\cdot\|}^+. \quad (3.9)$$

Let $\mathcal{M}_{\|\cdot\|}$ and $\mathcal{R}_{\|\cdot\|}$ denote the set of measures of the form $\nu = \nu_1 - \nu_2$ with $\nu_1, \nu_2 \in \mathcal{M}_{\|\cdot\|}^+$ or $\mathcal{R}_{\|\cdot\|}^+$ respectively. We often omit saying that both $\mathcal{R}_{\|\cdot\|}$ and $\|\cdot\|$ depend on the kernel u .

Let $\|\cdot\|$ be a proper norm for u . For $\nu_j \in \mathcal{R}_{\|\cdot\|}$, $j = 1, \dots, k$, let

$$M_t^{\nu_1, \dots, \nu_k} = \sum_{\pi \in \mathcal{T}_k^\odot} \int_{0 \leq r_1 \leq \dots \leq r_k \leq t} dL_{r_1}^{\nu_{\pi(1)}} \dots dL_{r_k}^{\nu_{\pi(k)}}. \quad (3.10)$$

We refer to $M_t^{\nu_1, \dots, \nu_k}$ as a multiple CAF. Let

$$Q^{x,y}(F) = \int_0^\infty Q_t^{x,y}(F \circ k_t) dt. \quad (3.11)$$

We have the following analogue of [9, Proposition 5] and [10, Lemma 2.1].

Lemma 3.1 *For any measures $\nu_j \in \mathcal{R}_{\|\cdot\|}^+$, $j = 1, \dots, k \geq 2$,*

$$\mu(M_\infty^{\nu_1, \dots, \nu_k}) = \int u(y_1, y_2) \dots u(y_{k-1}, y_k) u(y_k, y_1) \prod_{j=1}^k d\nu_j(y_j), \quad (3.12)$$

$$\begin{aligned} & Q^{x,y}(M_\infty^{\nu_1, \dots, \nu_k}) \\ &= \sum_{\pi \in \mathcal{T}_k^\odot} \int u(x, y_1) u(y_1, y_2) \dots u(y_{k-1}, y_k) u(y_k, y) \prod_{j=1}^k d\nu_{\pi(j)}(y_j), \end{aligned} \quad (3.13)$$

and if $\nu \in \mathcal{R}_{\|\cdot\|}^+$

$$\begin{aligned} & \mu(M_\infty^{\nu_1, \dots, \nu_k} L_\infty^\nu) \\ &= \sum_{i=1}^k \int \left(\prod_{j=1}^{i-1} u(y_j, y_{j+1}) \right) \\ & \quad u(y_i, x) u(x, y_{i+1}) \left(\prod_{j=i+1}^k u(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j) d\nu(x) \\ &= \int Q^{x,x}(M_\infty^{\nu_1, \dots, \nu_k}) d\nu(x), \end{aligned} \quad (3.14)$$

with $y_{k+1} = y_1$.

The proof of (3.12) follows that of [10, Lemma 2.1], noticing that the crucial step [10, (2.23)-(2.28)] used the fact that the set of permutations of $[1, k]$ is invariant under translation mod k . Since \mathcal{T}_k^\odot is invariant under translation mod k , the same proof will work here. The proof of (3.13), which is much easier, follows that of [10, Lemma 4.2]. The first equality in (3.14) follows from (3.12) and the fact that

$$M_\infty^{\nu_1, \dots, \nu_k} L_\infty^\nu = \sum_{i=1}^k M_\infty^{\nu_1, \dots, \nu_{i-1}, \nu, \nu_i, \dots, \nu_k}, \quad (3.15)$$

which is easily verified using $L_\infty^\nu = \int_0^\infty dL_t^\nu$. The second equality in (3.14) then follows by comparing with (3.13). \square

Let now $X_{(\epsilon)}$, $\epsilon \geq 0$, $X_{(0)} = X$, be a family of Markov processes with potential densities $u_{(\epsilon)}(x, y)$, and let $\mu_{(\epsilon)}$ denote the loop measure for $X_{(\epsilon)}$. Assume that we can use the same proper norm $\|\cdot\|$ for all $u_{(\epsilon)}$. Let

$$u'_{(0)}(x, y) = \frac{du_{(\epsilon)}(x, y)}{d\epsilon} \Big|_{\epsilon=0}, \quad (3.16)$$

and assume that $\|\cdot\|$ is also a proper norm for $u'_{(0)}$. Then using the last Lemma we have formally, that is, assuming we can justify interchanging derivative and integral in the second equality,

$$\begin{aligned} & \frac{d}{d\epsilon} \mu_{(\epsilon)}(M_\infty^{\nu_1, \dots, \nu_k}) \Big|_{\epsilon=0} \\ &= \frac{d}{d\epsilon} \int u_{(\epsilon)}(y_1, y_2) \cdots u_{(\epsilon)}(y_{k-1}, y_k) u_{(\epsilon)}(y_k, y_1) \prod_{j=1}^k d\nu_j(y_j) \Big|_{\epsilon=0} \\ &= \sum_{i=1}^k \int u(y_1, y_2) \cdots u'_{(0)}(y_i, y_{i+1}) \cdots u(y_k, y_1) \prod_{j=1}^k d\nu_j(y_j) \\ &= \sum_{\pi \in \mathcal{T}_k^\odot} \int u(y_1, y_2) \cdots u(y_{k-1}, y_k) u'_{(0)}(y_k, y_1) \prod_{j=1}^k d\nu_{\pi(j)}(y_j) \end{aligned} \quad (3.17)$$

where we have set $y_{k+1} = y_1$.

Assume now that for some distribution F on $S \times S$ we have

$$u'_{(0)}(y_k, y_1) = \int_{S \times S} u(y_k, x) F(x, y) u(y, y_1) dm(x) dm(y). \quad (3.18)$$

Let $\mathcal{A}_{\|\cdot\|}$ denote the space spanned by the multiple CAF's with $\nu_j \in \mathcal{R}_{\|\cdot\|}$. Note it is an algebra. Then comparing (3.17) and (3.13) we would obtain

$$\frac{d\mu_{(\epsilon)}(A)}{d\epsilon} \big|_{\epsilon=0} = \int_{S \times S} F(x, y) Q^{y,x}(A) dm(x) dm(y), \quad (3.19)$$

for all $A \in \mathcal{A}_{\|\cdot\|}$.

In the following sections we present specific examples where this heuristic approach is made rigorous.

4 Perturbation of Lévy processes

Let X be a transient Lévy process in R^d with characteristic exponent ψ so that, as distributions

$$u(x, y) = \int \frac{e^{i\lambda(y-x)}}{\psi(\lambda)} d\lambda. \quad (4.1)$$

In [10] we showed that $\|\cdot\|_{\psi,2}$ is a proper norm for u where

$$\|\nu\|_{\psi,2}^2 = \int \left(\frac{1}{|\psi|} * \frac{1}{|\psi|}(\lambda) \right) |\widehat{\nu}(\lambda)|^2 d\lambda. \quad (4.2)$$

Let κ be a Lévy characteristic exponent, so that the same is true for $\psi + \epsilon\kappa$, and let $X_{(\epsilon)}$ be the Lévy process with characteristic exponent $\psi + \epsilon\kappa$. We let $u_{(\epsilon)}(x, y)$ denote the potential of $X_{(\epsilon)}$ so that, as distributions

$$u_{(\epsilon)}(x, y) = \int \frac{e^{i\lambda(y-x)}}{\psi(\lambda) + \epsilon\kappa(\lambda)} d\lambda. \quad (4.3)$$

If we assume that

$$|\kappa(\lambda)| \leq C|\psi(\lambda)| \quad (4.4)$$

for some $C < \infty$, then for $\epsilon > 0$ sufficiently small

$$\|\nu\|_{\psi+\epsilon\kappa,2} \leq C'\|\nu\|_{\psi,2}, \quad (4.5)$$

for some $C' < \infty$. Thus $\|\cdot\|_{\psi,2}$ is a proper norm for $u_{(\epsilon)}$.

Let F be the distribution given by

$$F(x, y) = \int e^{i\lambda(y-x)} \kappa(\lambda) d\lambda, \quad (4.6)$$

and let $\widehat{Q}^{\lambda_1, \lambda_2}(A)$ denote the Fourier transform of $Q^{x,y}(A)$ in x, y .

Theorem 4.1 *If (4.4) holds, then*

$$\begin{aligned} \frac{d\mu_{(\epsilon)}(A)}{d\epsilon} \Big|_{\epsilon=0} &= - \int_{R^d \times R^d} Q^{y,x}(A) F(x,y) dm(x) dm(y) \\ &= - \int \widehat{Q}^{\lambda,-\lambda}(A) \kappa(\lambda) d\lambda. \end{aligned} \quad (4.7)$$

for all $A \in \mathcal{A}_{\|\cdot\|_{\psi,2}}$.

Proof of Theorem 4.1: It suffices to show that for any $\nu_1, \dots, \nu_k \in \mathcal{R}_{\|\cdot\|_{\psi,2}}^+$

$$\begin{aligned} & \frac{d}{d\epsilon} \int \prod_{j=1}^k u_{(\epsilon)}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \Big|_{\epsilon=0} \\ &= - \sum_{i=1}^k \int \left(\prod_{j=1}^{i-1} u(y_j, y_{j+1}) \right) \\ & \quad \left(\int_{R^d \times R^d} u(y_i, x) F(x, y) u(y, y_{i+1}) dm(x) dm(y) \right) \\ & \quad \left(\prod_{j=i+1}^k u(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j), \end{aligned} \quad (4.8)$$

with $y_{k+1} = y_1$.

Using (4.3) we see that

$$\begin{aligned} I(\epsilon) &=: \int \prod_{j=1}^k u_{(\epsilon)}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \\ &= \int \int \prod_{j=1}^k e^{i(y_{j+1}-y_j) \cdot \lambda_j} \frac{1}{\psi(\lambda_j) + \epsilon \kappa(\lambda_j)} d\lambda_j d\nu_j(y_j) \\ &= \int \left(\prod_{j=1}^k \int e^{-i(\lambda_j - \lambda_{j-1}) \cdot y_j} d\nu_j(y_j) \right) \prod_{j=1}^k \frac{1}{\psi(\lambda_j) + \epsilon \kappa(\lambda_j)} d\lambda_j \end{aligned} \quad (4.9)$$

where $\lambda_0 = \lambda_n$. We take the Fourier transforms of the $\{\nu_j\}$ to see that

$$\begin{aligned} & \int \prod_{j=1}^k u_{(\epsilon)}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \\ &= \int \prod_{j=1}^k \widehat{\nu}_j(\lambda_j - \lambda_{j-1}) \frac{1}{\psi(\lambda_j) + \epsilon \kappa(\lambda_j)} d\lambda_j. \end{aligned} \quad (4.10)$$

We have

$$\frac{1}{\psi(\lambda_j) + \epsilon \kappa(\lambda_j)} = \frac{1}{\psi(\lambda_j)} - \epsilon \frac{\kappa(\lambda_j)}{\psi^2(\lambda_j)} + \epsilon^2 \frac{\kappa^2(\lambda_j)}{\psi^2(\lambda_j)(\psi(\lambda_j) + \epsilon \kappa(\lambda_j))}. \quad (4.11)$$

Substituting this into (4.11) we can write the result as

$$I(\epsilon) = I(0) - \epsilon J + K(\epsilon), \quad (4.12)$$

where

$$J = \sum_{i=1}^k \int \left(\prod_{\substack{j=1 \\ j \neq i}}^k \frac{1}{\psi(\lambda_j)} \right) \frac{\kappa(\lambda_i)}{\psi^2(\lambda_i)} \prod_{j=1}^k \hat{\nu}_j(\lambda_j - \lambda_{j-1}) d\lambda_j, \quad (4.13)$$

and $K(\epsilon)$ is the sum of all remaining terms.

We now show that J is precisely the right hand side of (4.8), and that

$$|K(\epsilon)| = O(\epsilon^2), \quad (4.14)$$

which will complete the proof of our Proposition.

For the first point, using the relation between Fourier transforms and convolutions we have

$$\begin{aligned} \int_{R^d \times R^d} u(y_i, x) F(x, y) u(y, y_{i+1}) dm(x) dm(y) \\ = \int e^{i\lambda_i(y_{i+1} - y_i)} \frac{\kappa(\lambda_i)}{\psi^2(\lambda_i)} d\lambda_i. \end{aligned} \quad (4.15)$$

Using this in the right hand side of (4.8) and proceeding as in (4.9)-(4.10) we indeed obtain J . As for (4.14), $K(\epsilon)$ is the sum many terms each of which has a factor of ϵ^m for some $m \geq 2$. We need only show that the corresponding integrals are bounded uniformly in ϵ . For example, consider the term which arises when using the last term in (4.11) for all j . This term is ϵ^{2k} times

$$\tilde{K}(\epsilon) =: \int \prod_{j=1}^k \hat{\nu}_j(\lambda_j - \lambda_{j-1}) \frac{\kappa^2(\lambda_j)}{\psi^2(\lambda_j)(\psi(\lambda_j) + \epsilon \kappa(\lambda_j))} d\lambda_j. \quad (4.16)$$

By our assumption (4.4), for sufficiently small ϵ

$$\begin{aligned} |\tilde{K}(\epsilon)| &\leq C' \int \prod_{j=1}^k |\hat{\nu}_j(\lambda_j - \lambda_{j-1})| \frac{1}{|\psi(\lambda_j)|} d\lambda_j \\ &\leq C'' \prod_{j=1}^k \|\nu_j\|_{\psi, 2}, \end{aligned} \quad (4.17)$$

by repeated use of the Cauchy-Schwarz inequality, as in our proof in [10] that $\|\cdot\|_{\psi,2}$ is a proper norm for u . \square

5 Perturbation by multiplicative functionals

Let m_t be a continuous decreasing multiplicative functional of X , with $m_t \leq 1$ for all t and $m_\zeta = 0$. By [12, Theorem 61.5], there is a right process \tilde{X}_t with transition semigroup

$$\tilde{P}_t f(x) = P^x(f(X_t)m_t) = \int Q_t^{x,y}(m_t)f(y) dm(y), \quad (5.1)$$

where $Q_t^{x,y}$ are the bridge measures (2.1) for our original process X and the second equality is [6, (2.8)]. Thus \tilde{X} has transition densities

$$\tilde{p}_t(x, y) =: Q_t^{x,y}(m_t), \quad (5.2)$$

and using [6, (2.8)] once more we can verify that these satisfy the Chapman-Kolmogorov equations. It follows from the construction in [12, Theorem 61.5] that if X has cadlag paths so will \tilde{X} . Let $r_t(\omega)(s) = \omega(t-s)^-$, $0 \leq s \leq t$ be the time reversal mapping. If we set $\hat{m}_t = m_t \circ r_t$, then it is easy to check that \hat{m}_t is a multiplicative functional as above, see [2, p. 359]. If \hat{X} is the dual process for X as described in Section 2, with bridge measures $\hat{Q}_t^{x,y}$ then as above there exists a process Y with transition densities $\hat{Q}_t^{x,y}(\hat{m}_t)$. By [6, Corollary 1]

$$\hat{Q}_t^{x,y}(\hat{m}_t) = \hat{Q}_t^{x,y}(m_t \circ r_t) = Q_t^{y,x}(m_t), \quad (5.3)$$

which shows that Y is dual to \tilde{X} .

We now show that if $\tilde{Q}_t^{x,y}$ are the bridge measures for \tilde{X} , then

$$\tilde{Q}_t^{x,y}(F) = Q_t^{x,y}(F m_t), \text{ for } F \in \mathcal{F}_s, s < t. \quad (5.4)$$

To see this, using the fact that m_t is continuous and decreasing, we have for

$F \in \mathcal{F}_s, s < t$

$$\begin{aligned}
\tilde{Q}_t^{x,y}(F) &= \tilde{P}^x \left(F Q_{t-s}^{X_s,y}(m_{t-s}) \right) \\
&= \lim_{t^* \uparrow t} \tilde{P}^x \left(F Q_{t-s}^{X_s,y}(m_{t^*-s}) \right) \\
&= \lim_{t^* \uparrow t} \tilde{P}^x \left(F P^{X_s}(m_{t^*-s} p_{t-t^*}(X_{t^*-s}, y)) \right) \\
&= \lim_{t^* \uparrow t} P^x \left(F m_s P^{X_s}(m_{t^*-s} p_{t-t^*}(X_{t^*-s}, y)) \right) \\
&= \lim_{t^* \uparrow t} P^x \left(F m_s m_{t^*-s} \circ \theta_s p_{t-t^*}(X_{t^*}, y) \right) \\
&= \lim_{t^* \uparrow t} P^x \left(F m_{t^*} p_{t-t^*}(X_{t^*}, y) \right) \\
&= \lim_{t^* \uparrow t} Q_t^{x,y}(F m_{t^*}) = Q_t^{x,y}(F m_t),
\end{aligned} \tag{5.5}$$

which proves (5.4).

If A_t is a CAF, then $m_t = e^{-A_t}$ is a continuous decreasing multiplicative functional of X . Let $X_{(\epsilon)}$ denote the Markov process \tilde{X} with $m_t = e^{-\epsilon L_t^\nu}$. If $\mu_{(\epsilon)}$ denotes the loop measure for $X_{(\epsilon)}$, and μ the loop measure for X , it now follows from (2.2) and (5.4) that

$$\mu_{(\epsilon)}(A) = \mu(A e^{-\epsilon L_\infty^\nu}). \tag{5.6}$$

Theorem 5.1 *If $\nu \in \mathcal{R}_{\|\cdot\|}^+$, for some proper norm $\|\cdot\|$, then*

$$\frac{d\mu_{(\epsilon)}(A)}{d\epsilon} \Big|_{\epsilon=0} = -\mu(L_\infty^\nu A) = -\int_S Q^{x,x}(A) d\nu(x) \tag{5.7}$$

for all $A \in \mathcal{A}_{\|\cdot\|}$.

Proof: It suffices to prove this for A of the form $M_\infty^{\nu_1, \dots, \nu_k}$ with $\nu_j \in \mathcal{R}_{\|\cdot\|}^+$, $1 \leq j \leq k$. Since $0 \leq e^{-x} - 1 + x \leq x^2/2$ for $x \geq 0$, it follows from (5.6) that

$$|\mu_{(\epsilon)}(M_\infty^{\nu_1, \dots, \nu_k}) - \mu(M_\infty^{\nu_1, \dots, \nu_k}) - \epsilon \mu(L_\infty^\nu M_\infty^{\nu_1, \dots, \nu_k})| \leq \epsilon^2 \mu \left((L_\infty^\nu)^2 M_\infty^{\nu_1, \dots, \nu_k} \right). \tag{5.8}$$

$\mu \left((L_\infty^\nu)^2 A \right)$ is bounded by our assumption about the proper norm $\|\cdot\|$, so the first equality in (5.7) follows. The second equality is (3.14). \square

It follows from (5.2) that $X_{(\epsilon)}$ has potential densities

$$u_{(\epsilon)}(x, y) = \int_0^\infty Q_t^{x,y}(e^{-\epsilon L_t^\nu}) dt \leq u(x, y). \tag{5.9}$$

Then, using [6, Lemma 1]

$$\begin{aligned}
u'_{(0)}(x, y) &= - \int_0^\infty Q_t^{x,y}(L_t^\nu) dt \\
&= - \int_0^\infty \int_0^t \int p_s(x, z) p_{t-s}(z, y) d\nu(z) ds dt \\
&= - \int u(x, z) u(z, y) d\nu(z).
\end{aligned} \tag{5.10}$$

In view of this, Theorem 5.1 is another example of our heuristic formula (3.19), where the distribution F on $S \times S$ of (3.18) is $\delta(x - y) d\nu(x)$.

For use in the next section we now recast Theorem 5.1. For $\nu_j \in \mathcal{R}_{\|\cdot\|}^+$, $1 \leq j \leq k$ let

$$I(\epsilon) = \int \prod_{j=1}^k u_{(\epsilon)}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j). \tag{5.11}$$

Since $\int \prod_{j=1}^k u_{(\epsilon)}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) = \mu_{(\epsilon)}(M_\infty^{\nu_1, \dots, \nu_k})$ by (3.12), it follows from Theorem 5.1 and (3.14) that

$$\begin{aligned}
\lim_{\epsilon \rightarrow 0} \frac{I(\epsilon) - I(0)}{\epsilon} &= - \sum_{i=1}^k \int \left(\prod_{j=1}^{i-1} u(y_j, y_{j+1}) \right) u(y_i, x) u(x, y_{i+1}) \\
&\quad \left(\prod_{j=i+1}^k u(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j) d\nu(x).
\end{aligned} \tag{5.12}$$

In the next section we will also use the following observation. If $m_t = e^{-\epsilon \int_0^t c(s) ds}$ then $\widehat{m}_t = m_t \circ r_t = m_t$. Hence if \widehat{X} is the dual process for X as described in Section 2, it follows from (5.3) that Y , the dual process to $X_{(\epsilon)}$, is the process $(\widehat{X})_{(\epsilon)}$ obtained by perturbing \widehat{X} by the same multiplicative functional m_t . In particular

$$\widehat{u}_{(\epsilon)}(x, y) = u_{(\epsilon)}(y, x). \tag{5.13}$$

6 Perturbation by addition of jumps

Let $j(x, y)$ be a nonnegative $m \otimes m$ -integrable function on $S \times S$. We will assume that

$$c(x) =: \int j(x, y) m(dy) = \int j(y, x) m(dy). \tag{6.1}$$

$c(x)$ is integrable and we will assume moreover that it is bounded and strictly positive. Then

$$G(x, dy) =: \frac{1}{c(x)} j(x, y) m(dy) \quad (6.2)$$

is a probability kernel on $S \times \mathcal{B}(S)$. $c(x)$ will govern the rate at which we will add jumps to the process, which may depend on the position x of the process, and $G(x, dy)$ will describe the distribution of the jumps from position x .

In more detail, define the CAF

$$A_t = \int_0^t c(X_s) ds, \quad (6.3)$$

and let τ_t be the right continuous inverse of A_t . Let λ be an independent mean 1 exponential. We define a new process Y_t to be equal to X_t for $t < \tau_\lambda$, and then re-birthed at a random point independent of λ , distributed according to $G(X_{\tau_\lambda}, dy)$, with this process being iterated. We use $U_{c,G}$ to denote the potential operator of Y .

Let

$$\|\nu\|_{u^2, \infty} := |\nu|(S) \vee \sup_x \int (u^2(x, y) + u^2(y, x)) d|\nu|(y), \quad (6.4)$$

where $|\nu|$ is the total variation of the measure ν . This is a proper norm for u , see [10, (3.25)].

We use $\mu_{(\epsilon)}$ to denote the loop measure associated to Y , where we have replaced c by ϵc .

Theorem 6.1 *Assume that*

$$\sup_z \int u(z, y) dm(y) < \infty, \quad \sup_z \int u(y, z) dm(y) < \infty. \quad (6.5)$$

Then $\mu_{(\epsilon)}$ is well defined for ϵ small enough and

$$\frac{d\mu_{(\epsilon)}(A)}{d\epsilon} \Big|_{\epsilon=0} = \int_{S \times S} (Q^{y,x}(A) - Q^{x,x}(A)) c(x) G(x, dy) dm(x), \quad (6.6)$$

for all $A \in \mathcal{A}_{\|\cdot\|_{u^2, \infty}}$.

Before proving this theorem, we first show that $U_{\epsilon c, G}$ has densities $u_{\epsilon c, G}(x, y)$ for ϵ sufficiently small. Note first that by (6.1)-(6.2)

$$\int c(z) G(z, dy) dm(z) = \int j(z, y) dm(y) dm(z) = c(y) dm(y), \quad (6.7)$$

and since we assumed that c is bounded it follows from (6.5) that

$$\sup_x \int c(z) G(z, dy) u(y, x) dm(z) = \sup_x \int c(y) u(y, x) dm(y) < \infty. \quad (6.8)$$

Let $\lambda_1, \lambda_2, \dots$ be a sequence of independent mean 1 exponentials, and set $T_j = \sum_{i=1}^j \lambda_i$. Using the fact that $\tau_{t+u} = \tau_t + \tau_u \circ \theta_{\tau_t}$ we see that

$$\begin{aligned} W_{c,G,n} f(x) &= P^x \left(\int_{\tau(T_n)}^{\tau(T_{n+1})} f(Y_t) dt \right) \\ &= P^x \left(\left(\int_0^{\tau(\lambda_{n+1})} f(Y_t) dt \right) \circ \theta_{\tau(T_n)} \right) \\ &= P^x \left(P^{Y_{\tau(T_n)}} \left(\int_0^{\tau(\lambda_{n+1})} f(X_t) dt \right) \right) \\ &= P^x \left(\int G(Y_{\tau(T_n)}^-, dz) P^z \left(\int_0^{\tau(\lambda_{n+1})} f(X_t) dt \right) \right). \end{aligned} \quad (6.9)$$

We have

$$\begin{aligned} &P_\lambda^x \left(\int_0^{\tau_\lambda} f(X_t) dt \right) \\ &= P_\lambda^x \left(\int_0^\infty 1_{\{t < \tau_\lambda\}} f(X_t) dt \right) \\ &= P_\lambda^x \left(\int_0^\infty 1_{\{A_t < \lambda\}} f(X_t) dt \right) \\ &= P^x \left(\int_0^\infty e^{-A_t} f(X_t) dt \right). \end{aligned} \quad (6.10)$$

Hence setting

$$V_c f(x) = P^x \left(\int_0^\infty e^{-A_t} f(X_t) dt \right), \quad (6.11)$$

and writing $Gh(x) = \int_S G(z, dy) h(y)$ for any nonnegative or bounded function h , we have shown that

$$\begin{aligned} W_{c,G,n} f(x) &= P^x \left(\int G(Y_{\tau(T_n)}^-, dz) V_c f(z) \right) \\ &= P^x \left(G V_c f(Y_{\tau(T_n)}^-) \right). \end{aligned} \quad (6.12)$$

Using once again the fact that $\tau_{t+u} = \tau_t + \tau_u \circ \theta_{\tau_t}$ and the Markov property, we see that for any h

$$\begin{aligned}
P^x \left(h(Y_{\tau(T_n)}^-) \right) &= P^x \left(\left(h(Y_{\tau(\lambda_n)}^-) \right) \circ \theta_{\tau(T_{n-1})} \right) \\
&= P^x \left(P^{Y_{\tau(T_{n-1})}} \left(h(Y_{\tau(\lambda_n)}^-) \right) \right) \\
&= P^x \left(\int G(Y_{\tau(T_{n-1})}^-, dz) P^z \left(h(X_{\tau(\lambda_n)}^-) \right) \right).
\end{aligned} \tag{6.13}$$

Using the change of variables formula, [4, Chapter 6, (55.1)], and the fact that $X_t^- = X_t$ for a.e. t , we see that

$$\begin{aligned}
P_\lambda^z \left(h(X_{\tau_\lambda}^-) \right) = P_\lambda^z \left(h(X_{\tau_\lambda}) \right) &= E^x \left(\int_0^\infty e^{-t} h(X_{\tau_t}) dt \right) \\
&= P^z \left(\int_0^{A_\zeta} e^{-t} h(X_{\tau_t}) dt \right) \\
&= P^z \left(\int_0^\infty e^{-A_s} h(X_s) dA_s \right) \\
&= P^z \left(\int_0^\infty e^{-A_s} h(X_s) c(X_s) ds \right) \\
&= V_c(ch)(z).
\end{aligned} \tag{6.14}$$

Thus we can write (6.13) as

$$P^x \left(h(Y_{\tau(T_n)}^-) \right) = P^x \left(GV_c c h(Y_{\tau(T_{n-1})}^-) \right). \tag{6.15}$$

Iterating this we obtain

$$\begin{aligned}
P^x \left(h(Y_{\tau(T_n)}^-) \right) &= P^x \left(GV_c c h(Y_{\tau(T_{n-1})}^-) \right) \\
&= P^x \left((GV_c c)^2 h(Y_{\tau(T_{n-2})}^-) \right) \\
&= \dots \\
&= P^x \left((GV_c c)^{n-1} h(Y_{\tau(\lambda)}^-) \right) \\
&= V_c c (GV_c c)^{n-1} h(x),
\end{aligned} \tag{6.16}$$

where the last step used (6.14). Applying this to (6.12) we have that

$$W_{c,G,n} f(x) = V_c c (GV_c c)^{n-1} GV_c f(x) = V_c (cGV_c)^n f(x). \tag{6.17}$$

It follows from (5.9) that V_c has a density which we write as $v_c(x, y)$, and therefore $W_{c,G,n}$ has the density

$$\begin{aligned} w_{c,G,n}(x, y) = & \int v_c(x, z_1) G(z_1, dz_2) v_c(z_2, z_3) \cdots \\ & \cdots G(z_{2n-3}, dz_{2n-2}) v_c(z_{2n-2}, z_{2n-1}) G(z_{2n-1}, dz_{2n}) v_c(z_{2n}, y) \\ & \prod_{j=1}^n c(z_{2j-1}) dm(z_{2j-1}). \end{aligned} \quad (6.18)$$

By (6.5) it follows from (5.9) that

$$\sup_z \int v_c(z, y) dm(y) \leq M, \quad (6.19)$$

and thus by (6.18) that

$$\sup_z \int w_{c,G,n}(z, y) dm(y) \leq C^n M^{n+1}. \quad (6.20)$$

Replacing c by ϵc we have shown that for ϵ sufficiently small

$$U_{\epsilon c, G} f(x) = \sum_{n=0}^{\infty} \epsilon^n V_{\epsilon c} (c G V_{\epsilon c})^n f(x) \quad (6.21)$$

for all bounded measurable f . Hence $U_{\epsilon c, G}$ has a density

$$\begin{aligned} u_{\epsilon c, G}(x, y) = & \sum_{n=0}^{\infty} \epsilon^n \int v_{\epsilon c}(x, z_1) G(z_1, dz_2) v_{\epsilon c}(z_2, z_3) \cdots \\ & \cdots G(z_{2n-3}, dz_{2n-2}) v_{\epsilon c}(z_{2n-2}, z_{2n-1}) G(z_{2n-1}, dz_{2n}) v_{\epsilon c}(z_{2n}, y) \\ & \prod_{j=1}^n c(z_{2j-1}) dm(z_{2j-1}), \end{aligned} \quad (6.22)$$

with

$$\sup_z \int u_{\epsilon c, G}(z, y) dm(y) < \infty. \quad (6.23)$$

We can also write this as

$$\begin{aligned} u_{\epsilon c, G}(x, y) = & \sum_{n=0}^{\infty} \epsilon^n \int v_{\epsilon c}(x, z_1) j(z_1, z_2) v_{\epsilon c}(z_2, z_3) \cdots \\ & \cdots j(z_{2n-3}, z_{2n-2}) v_{\epsilon c}(z_{2n-2}, z_{2n-1}) j(z_{2n-1}, z_{2n}) v_{\epsilon c}(z_{2n}, y) \prod_{j=1}^{2n} dm(z_j). \end{aligned} \quad (6.24)$$

Similar expansions can be given for the semigroup which has therefore a density:

$$p_{\epsilon c, G, t}(x, y) = \sum_{n=0}^{\infty} \epsilon^n \int_{0 \leq t_1 \leq \dots \leq t_n \leq t} q_{t_1}(x, z_1) c(z_1) G(z_1, dz_2) q_{t_2-t_1}(z_2, z_3) \dots \\ \dots c(z_{2n-1}) G(z_{2n-1}, dz_{2n}) q_{t-t_n}(z_{2n}, y) \prod_{j=1}^n dm(z_{2j-1}) dt_j, \quad (6.25)$$

where q_t denotes the kernel of the semigroup associated with the process killed at rate ϵc .

To see this let

$$W_{c, G, n, t} f(x) =: P^x(f(Y_t); \tau(T_n) < t < \tau(T_{n+1})). \quad (6.26)$$

Using the fact that $\tau_{t+u} = \tau_t + \tau_u \circ \theta_{\tau_t}$ we have

$$\begin{aligned} &= P^x \left(f(Y_t); \tau(T_n) < t < \tau(T_{n+1}) \middle| \tau(T_n) \right) \\ &= P^x \left((f(Y_{t-\tau(T_n)}); t - \tau(T_n) < \tau(\lambda_{n+1})) \circ \theta_{\tau(T_n)} \middle| \tau(T_n) \right) \\ &= P^x \left(P^{Y_{\tau(T_n)}}(f(X_{t-\tau(T_n)}); t - \tau(T_n) < \tau(\lambda_{n+1})) \middle| \tau(T_n) \right) \\ &= P^x \left(\int G(Y_{\tau(T_n)}^-, dz) \right. \\ &\quad \left. P_{\lambda_{n+1}}^z(f(X_{t-\tau(T_n)}); t - \tau(T_n) < \tau(\lambda_{n+1})) \middle| \tau(T_n) \right). \end{aligned} \quad (6.27)$$

Conditional on $\tau(T_n)$ we have

$$\begin{aligned} &P_{\lambda}^x(f(X_{t-\tau(T_n)}); t - \tau(T_n) < \tau(\lambda)) \\ &= P_{\lambda}^x(f(X_{t-\tau(T_n)}); A_{t-\tau(T_n)} < \lambda) \\ &= P^x(e^{-A_{t-\tau(T_n)}} f(X_{t-\tau(T_n)})). \end{aligned} \quad (6.28)$$

Hence setting

$$P_{c, t} f(x) = P^x(e^{-A_t} f(X_t)), \quad (6.29)$$

we have shown that

$$\begin{aligned} W_{c, G, n, t} f(x) &= P^x \left(\int G(Y_{\tau(T_n)}^-, dz) P_{c, t-\tau(T_n)} f(z) \right) \\ &= P^x \left(G P_{c, t-\tau(T_n)} f(Y_{\tau(T_n)}^-) \right). \end{aligned} \quad (6.30)$$

Using once again the fact that $\tau_{t+u} = \tau_t + \tau_u \circ \theta_{\tau_t}$ and the strong Markov property, we see that for any h

$$\begin{aligned}
& P^x \left(h(Y_{\tau(T_n)}^-, t - \tau(T_n)) \middle| \tau(T_{n-1}) \right) \\
&= P^x \left(\left(h(Y_{\tau(\lambda_n)}^-, t - \tau(T_{n-1}) - \tau(\lambda_n)) \right) \circ \theta_{\tau(T_{n-1})} \middle| \tau(T_{n-1}) \right) \\
&= P^x \left(P^{Y_{\tau(T_{n-1})}} \left(h(Y_{\tau(\lambda_n)}^-, t - \tau(T_{n-1}) - \tau(\lambda_n)) \right) \middle| \tau(T_{n-1}) \right) \\
&= P^x \left(\int G(Y_{\tau(T_{n-1})}^-, dz) P^z \left(h(X_{\tau(\lambda_n)}^-, t - \tau(T_{n-1}) - \tau(\lambda_n)) \right) \middle| \tau(T_{n-1}) \right).
\end{aligned} \tag{6.31}$$

Using the change of variables formula, [4, Chapter 6, (55.1)], and the fact that $X_t^- = X_t$ for a.e. t , we see that conditionally on $\tau(T_{n-1})$

$$\begin{aligned}
& P_{\lambda_n}^z \left(h(X_{\tau(\lambda_n)}^-, t - \tau(T_{n-1}) - \tau(\lambda_n)) \right) \\
&= P_{\lambda_n}^z \left(h(X_{\tau(\lambda_n)}, t - \tau(T_{n-1}) - \tau(\lambda_n)) \right) \\
&= E^x \left(\int_0^\infty e^{-r} h(X_{\tau_r}, t - \tau(T_{n-1}) - \tau_r) dr \right) \\
&= P^z \left(\int_0^{A_\zeta} e^{-r} h(X_{\tau_r}, t - \tau(T_{n-1}) - \tau_r) dr \right) \\
&= P^z \left(\int_0^\infty e^{-A_s} h(X_s, t - \tau(T_{n-1}) - s) dA_s \right) \\
&= \int_0^\infty P^z (e^{-A_s} h(X_s, t - \tau(T_{n-1}) - s) c(X_s)) ds \\
&= \int_0^\infty P_{c,s}(ch(\cdot, t - s - \tau(T_{n-1}))) (z) ds.
\end{aligned} \tag{6.32}$$

Combining this with (6.31) we have

$$\begin{aligned}
& P^x \left(h(Y_{\tau(T_n)}^-, t - \tau(T_n)) \right) \\
&= \int P^x \left(GP_{c,s_n} ch(Y_{\tau(T_{n-1})}^-, t - s_n - \tau(T_{n-1})) \right) ds_n.
\end{aligned} \tag{6.33}$$

Iterating this we obtain, conditionally on $\tau(T_{n-1})$, then $\tau(T_{n-2}), \dots$

$$\begin{aligned}
& P^x \left(h(Y_{\tau(T_n)}^-), t - \tau(T_n) \right) \\
&= \int P^x \left(GP_{c,s_n} ch(Y_{\tau(T_{n-1})}^-), t - s_n - \tau(T_{n-1}) \right) ds_n \\
&= \int P^x \left(GP_{c,s_{n-1}} cGP_{c,s_n} ch(Y_{\tau(T_{n-2})}^-), t - s_n - s_{n-1} - \tau(T_{n-2}) \right) ds_{n-1} ds_n \\
&= \dots \\
&= \int P^x \left(GP_{c,s_2} c \dots GP_{c,s_n} ch(Y_{\tau(\lambda)}^-), t - \sum_{j=2}^n s_j - \tau(\lambda) \right) \prod_{j=2}^n ds_j \\
&= \int P_{c,s_1} cGP_{c,s_2} c \dots GP_{c,s_n} ch(x, t - \sum_{j=1}^n s_j) \prod_{j=1}^n ds_j
\end{aligned} \tag{6.34}$$

where the last step used (6.32). Applying this to (6.30) we have that

$$W_{c,G,n,t} f(x) = \int P_{c,s_1} cGP_{c,s_2} c \dots GP_{c,s_n} cGP_{c,t-\sum_{j=1}^n s_j} f(x) \prod_{j=1}^n ds_j. \tag{6.35}$$

(6.25) then follows as in the proof of (6.22).

Let

$$\hat{G}(x, dy) = \frac{1}{c(x)} j(y, x) m(dy). \tag{6.36}$$

By (6.1), $\hat{G}(x, dy)$ is a probability kernel on $S \times \mathcal{B}(S)$. We now add jumps to the dual process \widehat{X} , where $c(x)$ will again govern the rate of jumps but we use $\hat{G}(x, dy)$ to describe the distribution of the jumps from position x . Performing the same calculation as before, but with the dual process and using (6.24) and (5.13), we see that the two processes obtained by adding jumps have dual potential kernels:

$$\hat{u}_{\epsilon c, \hat{G}}(x, y) = u_{\epsilon c, G}(y, x) \tag{6.37}$$

The same will be true for the associated resolvents and semigroups. The duality assumptions are verified and we can therefore define a loop measure associated with these processes.

We use $\mu_{(\epsilon)}$ to denote the loop measure associated to Y , where we have replaced c by ϵc .

The next Lemma is needed for the proof of Theorem 6.1.

Lemma 6.1 *Assume (6.5) (which implies (6.8)). Then for any positive measure ν and ϵ sufficiently small.*

$$\sup_z \int u_{\epsilon c, G}(z, y) d\nu(y) \leq C \sup_z \int u(z, y) d\nu(y), \quad (6.38)$$

and

$$\|\nu\|_{u_{\epsilon c, G}^2, \infty} \leq C \|\nu\|_{u^2, \infty}, \quad (6.39)$$

Proof of Lemma 6.1: (6.38) follows immediately from (6.22) and (5.9). It also follows from (6.22) that for a positive measure ν

$$\begin{aligned} & \int u_{\epsilon c, G}^2(x, y) d\nu(y) \\ &= \sum_{m, n=0}^{\infty} \epsilon^{m+n} \int V_{\epsilon c} (cGV_{\epsilon c})^m(x, y) V_{\epsilon c} (cGV_{\epsilon c})^n(x, y) d\nu(y) \\ &= \sum_{m, n=0}^{\infty} \epsilon^{m+n} \int V_{\epsilon c} (cGV_{\epsilon c})^{m-1} cG(x, dz_1) V_{\epsilon c} (cGV_{\epsilon c})^{n-1} cG(x, dz_2) \\ & \quad \left(\int v_{\epsilon c}(z_1, y) v_{\epsilon c}(z_2, y) d\nu(y) \right). \end{aligned} \quad (6.40)$$

Hence for ϵ small enough

$$\begin{aligned} & \sup_x \int u_{\epsilon c, G}^2(x, y) d\nu(y) \\ & \leq \sum_{m, n=0}^{\infty} \epsilon^{m+n} \sup_x \int V_{\epsilon c} (cGV_{\epsilon c})^{m-1} cG(x, dz_1) V_{\epsilon c} (cGV_{\epsilon c})^{n-1} cG(x, dz_2) \\ & \quad \sup_{z_1, z_2} \left(\int v_{\epsilon c}(z_1, y) v_{\epsilon c}(z_2, y) d\nu(y) \right) \\ & \leq C \|\nu\|_{u^2, \infty}. \end{aligned} \quad (6.41)$$

Similarly

$$\begin{aligned} & \int u_{\epsilon c, G}^2(y, x) d\nu(y) \\ &= \sum_{m, n=0}^{\infty} \epsilon^{m+n} \int V_{\epsilon c} (cGV_{\epsilon c})^m(y, x) V_{\epsilon c} (cGV_{\epsilon c})^n(y, x) d\nu(y) \\ &= \sum_{m, n=0}^{\infty} \epsilon^{m+n} \int (cGV_{\epsilon c})^m(z_1, x) (cGV_{\epsilon c})^n(z_2, x) dm(z_1) dm(z_2) \\ & \quad \left(\int v_{\epsilon c}(y, z_1) v_{\epsilon c}(y, z_2) d\nu(y) \right). \end{aligned} \quad (6.42)$$

Using (6.8) it follows that for ϵ small enough

$$\begin{aligned}
& \sup_x \int u_{\epsilon c, G}^2(x, y) d\nu(y) \\
& \leq \sum_{m, n=0}^{\infty} \epsilon^{m+n} \sup_x \int (cGV_{\epsilon c})^m(z_1, x) (cGV_{\epsilon c})^n(z_2, x) dm(z_1) dm(z_2) \\
& \quad \sup_{z_1, z_2} \left(\int v_{\epsilon c}(y, z_1) v_{\epsilon c}(y, z_2) d\nu(y) \right) \\
& \leq C \|\nu\|_{u^2, \infty}.
\end{aligned} \tag{6.43}$$

□

It follows from [10, Lemma 3.3] that $\|\nu\|_{u^2, \infty}$ is a proper norm for $u_{\epsilon c, G, n}$. Theorem 6.1 follows from the next Lemma.

Lemma 6.2 *Under the assumptions of Lemma 6.1, for any $\nu_1, \dots, \nu_k \in \mathcal{R}_{\|\cdot\|_{u^2, \infty}}^+$*

$$\begin{aligned}
& \frac{d}{d\epsilon} \int \prod_{j=1}^k u_{\epsilon c, G}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \Big|_{\epsilon=0} \\
& = \sum_{i=1}^k \int \left(\prod_{j=1}^{i-1} u(y_j, y_{j+1}) \right) \left(- \int u(y_i, x) u(x, y_{i+1}) c(x) dm(x) \right. \\
& \quad \left. + \int_{S \times S} u(y_i, z_1) c(z_1) G(z_1, dz_2) u(z_2, y_{i+1}) dm(z_1) \right) \\
& \quad \left(\prod_{j=i+1}^k u(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j),
\end{aligned} \tag{6.44}$$

with $y_{k+1} = y_1$.

Proof of Lemma 6.2: Set

$$\mathcal{I}(\epsilon) = \int \prod_{j=1}^k u_{\epsilon c, G}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \tag{6.45}$$

with $y_{k+1} = y_1$.

We can write (6.22) as

$$\begin{aligned}
& u_{\epsilon c, G}(x, y) \\
&= v_{\epsilon c}(x, y) + \epsilon \int v_{\epsilon c}(x, z_1) c(z_1) G(z_1, dz_2) v_{\epsilon c}(z_2, y) dm(y_1) \\
&+ \epsilon^2 \int v_{\epsilon c}(x, z_1) c(z_1) G(z_1, dz_2) v_{\epsilon c}(z_2, y) c(z_2) G(z_2, dz_3) u_{\epsilon c, G}(z_3, y) dm(y_1) dm(y_3) \\
&= v_{\epsilon c}(x, y) + \epsilon V_{\epsilon c} c G V_{\epsilon c}(x, y) + \epsilon^2 V_{\epsilon c} c G V_{\epsilon c} c G U_{\epsilon c, G}(x, y),
\end{aligned} \tag{6.46}$$

with operator notation.

We substitute this in (6.45) and collect terms to obtain

$$\mathcal{I}(\epsilon) = I(\epsilon) + \epsilon \sum_{i=1}^k \mathcal{J}_i(\epsilon) + \mathcal{K}(\epsilon), \tag{6.47}$$

where

$$I(\epsilon) = \int \prod_{j=1}^k v_{\epsilon c}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \tag{6.48}$$

$$\begin{aligned}
& \mathcal{J}_i(\epsilon) \\
&= \int \left(\prod_{j=1}^{i-1} v_{\epsilon c}(y_j, y_{j+1}) \right) V_{\epsilon c} c G V_{\epsilon c}(y_i, y_{i+1}) \left(\prod_{j=i+1}^k v_{\epsilon c}(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j),
\end{aligned} \tag{6.49}$$

and $\mathcal{K}(\epsilon)$ represents all the remaining terms. Noting that $\mathcal{I}(0) = I(0)$, we can write (6.47) as

$$\mathcal{I}(\epsilon) - \mathcal{I}(0) = I(\epsilon) - I(0) + \epsilon \sum_{i=1}^k \mathcal{J}_i(\epsilon) + \mathcal{K}(\epsilon). \tag{6.50}$$

Note also that $I(\epsilon)$ of (6.48) is a special case of the $I(\epsilon)$ of (5.11) with $\nu(dx) = c(x) dm(x)$. Hence by (5.12)

$$\begin{aligned}
& \lim_{\epsilon \rightarrow 0} \frac{I(\epsilon) - I(0)}{\epsilon} \\
&= - \sum_{i=1}^k \int \left(\prod_{j=1}^{i-1} u(y_j, y_{j+1}) \right) \left(\int u(y_i, x) u(x, y_{i+1}) c(x) dm(x) \right) \\
&\quad \left(\prod_{j=i+1}^k u(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j).
\end{aligned} \tag{6.51}$$

Since $v_{\epsilon c}(x, y) \uparrow u(x, y)$ as $\epsilon \downarrow 0$, it follows by the Monotone Convergence Theorem that

$$\begin{aligned} & \lim_{\epsilon \rightarrow 0} \mathcal{J}_i(\epsilon) \\ &= \int \left(\prod_{j=1}^{i-1} u(y_j, y_{j+1}) \right) \\ & \quad \left(\int_{S \times S} u(y_i, z_1) c(z_1) G(z_1, dz_2) u(z_2, y_{i+1}) dm(z_1) \right) \\ & \quad \left(\prod_{j=i+1}^k u(y_j, y_{j+1}) \right) \prod_{j=1}^k d\nu_j(y_j). \end{aligned} \tag{6.52}$$

To complete the proof of our Lemma it remains to show that

$$\lim_{\epsilon \rightarrow 0} \frac{\mathcal{K}(\epsilon)}{\epsilon} = 0. \tag{6.53}$$

However every term in $\mathcal{K}(\epsilon)$ comes with a pre-factor of ϵ^m for some $m \geq 2$, so we need only bound the integrals uniformly in ϵ . For this we will use Lemma 6.1. We illustrate this with the most complicated term, which has the pre-factor ϵ^{2k} :

$$\begin{aligned} & \int \prod_{j=1}^k V_{\epsilon c} c G V_{\epsilon c} c G U_{\epsilon c, G}(y_j, y_{j+1}) \prod_{j=1}^k d\nu_j(y_j) \\ &= \int v_{\epsilon c}(y_1, x_1) c G V_{\epsilon c} c G U_{\epsilon c, G}(x_1, y_2) \cdots \\ & \quad v_{\epsilon c}(y_{k-1}, x_{k-1}) c G V_{\epsilon c} c G U_{\epsilon c, G}(x_{k-1}, y_k) \\ & \quad v_{\epsilon c}(y_k, x_k) c G V_{\epsilon c} c G U_{\epsilon c, G}(x_k, y_1) \prod_{j=1}^k d\nu_j(y_j) dm(x_j) \\ &= \int c G V_{\epsilon c} c G U_{\epsilon c, G}(x_1, y_2) v_{\epsilon c}(y_2, x_2) \cdots \\ & \quad c G V_{\epsilon c} c G U_{\epsilon c, G}(x_{k-1}, y_k) v_{\epsilon c}(y_k, x_k) \\ & \quad c G V_{\epsilon c} c G U_{\epsilon c, G}(x_k, y_1) v_{\epsilon c}(y_1, x_1) \prod_{j=1}^k d\nu_j(y_j) dm(x_j), \end{aligned} \tag{6.54}$$

where the last step is just a rearrangement. We can rewrite this as

$$\begin{aligned} & \int \left(\int cGV_{\epsilon c} cGU_{\epsilon c, G}(x_1, y_2) v_{\epsilon c}(y_2, x_2) d\nu_2(y_2) \right) \cdots \\ & \left(\int cGV_{\epsilon c} cGU_{\epsilon c, G}(x_{k-1}, y_k) v_{\epsilon c}(y_k, x_k) d\nu_k(y_k) \right) \\ & \left(\int cGV_{\epsilon c} cGU_{\epsilon c, G}(x_k, y_1) v_{\epsilon c}(y_1, x_1) d\nu_1(y_1) \right) \prod_{j=1}^k dm(x_j). \end{aligned} \quad (6.55)$$

Then by Lemma 6.1, and the fact that $v_{\epsilon c} \leq u$

$$\begin{aligned} & \int cGV_{\epsilon c} cGU_{\epsilon c, G}(x_1, y_2) v_{\epsilon c}(y_2, x_2) d\nu_2(y_2) \\ & = \int c(x_1) G(x_1, dz_1) v_{\epsilon c}(z_1, z_2) c(z_2) G(z_2, dz_3) u_{\epsilon c, G}(z_3, y_2) v_{\epsilon c}(y_2, x_2) d\nu_2(y_2) dm(z_2) \\ & \leq \int c(x_1) G(x_1, dz_1) v_{\epsilon c}(z_1, z_2) c(z_2) G(z_2, dz_3) dm(z_2) \\ & \quad \sup_{z_3, x_2} \int u_{\epsilon c, G}(z_3, y_2) v_{\epsilon c}(y_2, x_2) d\nu_2(y_2) \\ & \leq C \|\nu_2\|_{u^2, \infty} \int c(x_1) G(x_1, dz_1) v_{\epsilon c}(z_1, z_2) c(z_2) G(z_2, dz_3) dm(z_2). \end{aligned} \quad (6.56)$$

Using (6.19), the fact that c is bounded and the fact that $G(\cdot, dz)$ is a probability density

$$\int c(x_1) G(x_1, dz_1) v_{\epsilon c}(z_1, z_2) c(z_2) G(z_2, dz_3) dm(z_2) \leq C c(x_1). \quad (6.58)$$

Thus (6.55) is bounded independently of ϵ by

$$C \int \prod_{j=1}^k c(x_j) dm(x_j) < \infty \quad (6.59)$$

since $c(x)$ is integrable.

The other terms of $\mathcal{K}(\epsilon)$ can be bounded similarly. \square

References

1. R. Blumenthal and R. Gettoor, *Markov Processes and Potential Theory*, Academic Press, New York, 1968. 2
2. K. L. Chung and J. Walsh, *Markov Processes, Brownian Motion and Time Symmetry*, second edition, Springer, New York, 2005. 5
3. Y. Chang, Multi-occupation field generates the Borel sigma-field of loops. arXiv:1309.1558 [math.PR] 2
4. C. Dellacherie, and P.-A. Meyer, (1982). *Probabilities and Potential B*. North Holland Publishing Company, Amsterdam. 6, 6
5. P. J. Fitzsimmons and J. Rosen, Markovian loop soups: permanental processes and isomorphism theorems, arxiv.org/pdf/1211.5163.pdf 2
6. P. Fitzsimmons, J. Pitman, and M. Yor, Markovian bridges: construction, Palm interpretation, and splicing. *Seminar on Stochastic Processes*, 1992, E. Cinlar and K.L. Chung and M.J. Sharpe editors, 101-134, Birkhuser, Boston (1993). 2, 2, 5, 5, 5
7. M. Fukushima, Y. Oshima and M. Takeda, *Dirichlet Forms and Symmetric Markov Processes*. Walter de Gruyter, New York, 1994. 1
8. K. Gawedski, *Lectures on conformal field theory*, IAS, Princeton, (2009). www.math.ias.edu/QFT/fall/ 1
9. Y. Le Jan, *Markov paths, loops and fields*. École d'Été de Probabilités de Saint-Flour XXXVIII - 2008. Lecture Notes in Mathematics 2026. (2011) Springer-Verlag, Berlin-Heidelberg. 1, 2, 2, 3
10. Y. Le Jan, M. B. Marcus and J. Rosen, Permanental fields, loop soups and continuous additive functionals, arxiv.org/pdf/1209.1804.pdf, Version 1. 3, 3, 3, 3, 4, 4, 6, 6
11. Y. Le Jan, M. B. Marcus and J. Rosen, Intersection local times, loop soups and permanental Wick powers, arxiv.org/pdf/1308.2701.pdf 2
12. M. Sharpe, *General theory of Markov processes*, Acad. Press, New York, 1988. 2, 5, 5

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